# IMBEDDING OF INFINITE DIMENSIONAL DISTRIBUTIONS INTO 

# SIMPLIFIED COLOMBEAU TYPE ALGEBRAS 

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#### Abstract

Using a d-dimensional Gaussian probability space, a simplified Colombeau-type algebra is constructed containing the Meyer-Watanabe distributions. A parallel construction starting by a certain Gelfand triplet also includes the Hida distributions. As an application a new interpretation of the Feynman integrand as a generalized function in our sense is proposed.


Keywords. Asymptotic Colombeau extension, Meyer-Watanabe distributions, Hida distributions, Donsker's delta function, sharp uniform topology, Feynman integrand.

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## 1 Introduction

It is well-known that the Colombeau algebra developed in 1984 contained the distribution spaces $\mathcal{D}^{\prime}$ and $\mathcal{S}^{\prime}$ ([1], [2], [3], [4]) and warded off the consequences of the Schwartz's Impossibility Result and Stability Paradox in an optimal manner, i.e. the concessions given were kept at a minimum. In return there was a well-defined rule of multiplication for distributions. In a simplified version of the Colombeau algebra, the functional index set which is an essential feature in Colombeau's theory is avoided at the expense of non-existence of a privileged inclusion of $\mathcal{D}^{\prime}$ distributions. However infinite number representatives are related within equivalence relations.

Since then various aspects, extensions and ramifications of the Colombeau's theory have been investigated including its random version, use in the solution of non-linear p.d.e. and relation to the non-standard analysis and ultradistributions, ([5], [6], [7], [8],[9], [10], [11], [12], [13], [14]).
However there have not been noticable attempts to construct associative, commutative algebras containing infinite dimensional distributions with a possible exception of [15] where Hida distributions are imbedded into an algebra via chaos expansions.

In Section 2, inspired by the Colombeau a-extension proposed by Delcroix \& Scarpalezos ([16]) we start by a ddimensional Gaussian space and construct an asymptotic, polynomial scale Colombeau extension $\mathcal{H}_{s}$. We show that the Meyer-Watanabe distributions $\mathbb{D}^{-\infty}$ are included in the algebra $\mathcal{H}_{s}$.
In Section 3 we consider a particular Gelfand triplet $E \subset H \subset E^{*}$ and the related Hida distribution. Using the Gaussian measure $\mu$ on $E^{*}$ provided by the Minlos theorem we construct a parallel of the algebra $\mathcal{H}_{s}$ which also contains the Hida distributions. Thus in $\mathcal{H}_{s}$ the product of Hida distributions is well-defined in comparison to the indirect and somewhat artificial Wick product via the inverse $S$-transform.
After extending the results to the complex case in subsection 3.2, we give as an application, a new interpretation of the Feynman integrand as a distribution in our sense in 3.3.

The present article is based on and the extension of the results taking place in unpublished talks given by the author in international conferences [17], [ 18].

## 2 Inclusion of Meyer-Watanabe Distributions

### 2.1 Basic Framework and Notation

$\Omega=C_{0}([0,1])$ : The set of all $\mathbb{R}^{d}$-valued continuous functions on $[0,1]$, null at zero with the sup norm.
$\mu$ : the Wiener measure on $(\Omega, \mathcal{B}(\Omega))$.
For $t \in[0,1], \omega \in \Omega, W_{t}(\omega) \doteq \omega(t)$ is an $\mathbb{R}^{d}$-valued Wiener functional and $W(t, \omega)$ is a $d$-dimensional Wiener process.
$\mathcal{F}_{t}: t \in[0,1]:$ the natural filtration generated by the process. $\mathcal{F} \equiv \mathcal{F}_{\infty}$ is the $\mu$-completion of the sigma algebra $\mathcal{B}(\Omega)$, so that $(\Omega, \mathcal{F}, \mu)$ is our Wiener space.
$H$ : the Cameron-Martin (C-M) space, is the Hilbert space formed by all $h \in \Omega$ such that each component of $\left(h(t)=\left(h_{1}(t), h_{2}(t), \cdots, h_{d}(t)\right)\right.$ is absolutely continuous and has square integrable derivatives.
The inner product in $H$ : for $g, h \in H$

$$
\begin{equation*}
(g, h)_{H}=\int_{0}^{1} \sum_{1}^{d} \dot{g}_{i}(t) \dot{h_{i}}(t) d t \text { and the Hilbertian norm }|h|_{H}^{2}=\int_{0}^{1}|\dot{h}(t)|^{2} d t \tag{2.1.1}
\end{equation*}
$$

An $\mathcal{F}$-measurable fuction $F: \Omega \rightarrow \mathbb{R}$ is a Wiener functional. Polynomial (resp. smooth) functionals are denoted by $\mathcal{P}\left(\right.$ resp. $\left.S_{M}\right)$ We have $\mathcal{P} \subset S_{M} \subset L^{p} \equiv L^{p}(\Omega, \mathcal{F}, \mu)$, also $\mathcal{P}$ is dense in $L^{p},(1 \leq p<\infty)$.
Ocassionally we use $E$-valued functionals $F: \Omega \rightarrow E$, where $E$ is a separable Hilbert space. In this case the symbols are $\mathcal{P}(E) \subset S_{M}(E) \subset L^{p}(E)$.
$D_{h}$ denotes the perturbation of Wiener functionals in the direction of members of the C-M space, i.e. the weak derivative defined by the well-known formula $D_{h} F(\omega)=\lim _{\lambda \rightarrow 0} \frac{F(\omega+\lambda h)-F(\omega)}{h} ; h \in H, F \in S_{M}$.
The gradient operator $D$ is the closable operator from $L^{p}$ into $L^{p}(H)$ determined by

$$
\begin{equation*}
(D F, h)_{H}=D_{h} F ; D F \in S_{M}(H) \tag{2.1.2}
\end{equation*}
$$

For $E$-valued functionals, the counterpart of (2.1.2) is

$$
\begin{equation*}
(D F, h \otimes e)_{H \otimes E}=\lim _{\lambda \rightarrow 0} \frac{1}{\lambda}(F(\omega+\lambda h)-F(\omega), e)_{E} ; e \in E, D F \in S_{M}(H \otimes E) \tag{2.1.3}
\end{equation*}
$$

The chain rule for the gradients is

$$
\begin{equation*}
D\left(\phi\left(F_{1}, F_{2}, \cdots F_{n}\right)=\sum_{j=1}^{n} \partial_{j} \phi\left(F_{1}, F_{2}, \cdots F_{n}\right) D F_{i}\right. \tag{2.1.4}
\end{equation*}
$$

where $F_{i} \in S_{M},(i=1, \cdots, n)$ and $\phi$ is a tempered $\mathcal{C}^{\infty}$-function on $\mathbb{R}^{n}$, (in fact it suffices that $\phi$ is smooth, c.f. ([19], II, (2.26))

The Ornstein-Uhlenbeck operator is $\mathcal{L}=-\delta D$ where $\delta$ is the divergence operator, i.e. the adjoint of $D$.
The norms defined on polynomial functionals: $\|F\|_{s, p}=\left\|(I-\mathcal{L})^{s / 2} F\right\|_{p}, s \in \mathbb{R}, 1<p<\infty$ have the properties of monotonicity and consistency. It is known that for $s=k \in \mathbb{N}$

$$
\begin{equation*}
\|F\|_{k, p}=\left(\|F\|_{L^{p}}^{p}+\sum_{j=1}^{k}\left\|D^{j} F\right\|_{L^{p}\left(\mu ; H^{\otimes j)}\right.}^{p}\right)^{\frac{1}{p}} \tag{2.1.5}
\end{equation*}
$$

where $D^{j} F=D\left(D^{j-1} F\right), \in S_{M}\left(H^{\otimes j}\right)$, (the tensor products are symmetrized).
We denote by $\mathbb{D}_{s}^{p}$, the Banach space which is the completion of $S_{M}$ with respect to the norm $\|.\|_{s, p}$ and

$$
\begin{equation*}
\mathbb{D}^{\infty}=\bigcap_{s>0} \bigcap_{1<p<\infty} \mathbb{D}_{s}^{p}, \quad \mathbb{D}^{-\infty}=\bigcup_{s>0} \bigcup_{1<p<\infty} \mathbb{D}_{-s}^{p} \tag{2.1.6}
\end{equation*}
$$

are Meyer-Watanabe testing functional space and its dual space (Meyer-Watanabe distributions) respectively.
$\mathbb{D}^{\infty}$ is a complete countably normed space (Frechet space) and also a topological algebra. By the monotonicity of $\|\cdot\|_{s, p}$ norms we also have

$$
\begin{equation*}
\mathbb{D}^{\infty}=\bigcap_{k \in \mathbb{N}} \bigcap_{1<p<\infty} \mathbb{D}_{k}^{p} \tag{2.1.7}
\end{equation*}
$$

Natural generalizations to $E$-valued functionals would be

$$
\mathbb{D}^{\infty}(E)=\bigcap_{s>0} \bigcap_{1<p<\infty} \mathbb{D}_{s}^{p}(E), \mathbb{D}^{-\infty}(E)=\bigcup_{s>0} \bigcup_{1<p<\infty} \mathbb{D}_{-s}^{p}(E)
$$

In many respects the pair $\left(\mathbb{D}^{\infty}, \mathbb{D}^{-\infty}\right)$ behaves like $\left(\mathcal{D}, \mathcal{D}^{\prime}\right)$ of the Schwartz theory.

### 2.2 Construction of the Algebra $\mathcal{H}_{s}$

Let $\chi$ be the set of all functions from $(0,1]$ into $\mathbb{D}^{\infty}$. A member of $\chi$ will then be denoted by $\left(f_{\epsilon}\right)_{\epsilon} \in \chi \equiv\left(\mathbb{D}^{\infty}\right)^{(0,1]}$.

Definition 2.2.1. $\left(f_{\epsilon}\right)_{\epsilon} \in \chi$ is called moderate if given $(k, p), k \in \mathbb{N}, 1<p<\infty, \exists N \in \mathbb{N}$ such that $\left\|f_{\epsilon}\right\|_{k, p}=O\left(\epsilon^{-N}\right)$ as $\epsilon \rightarrow 0$.
The set of all moderate elements in $\chi$ is denoted by $\mathcal{M}_{s}=(\Omega, \mathcal{F}, \mu ; H)$.
Definition 2.2.2. $\left(f_{\epsilon}\right)_{\epsilon} \in \chi$ is called negligible (or null), if given $(k, p), k \in \mathbb{N},, 1<p<\infty, \exists N \in \mathbb{N}$ such that $\left\|f_{\epsilon}\right\|_{k, p}=O\left(\epsilon^{m-N}\right), \forall m \geq N$ as $\epsilon \rightarrow 0$.
The set of all negligible members of $\chi$ is denoted by $\mathcal{J}_{s}(\Omega, \mathcal{F}, \mu ; H)$.

Definition 2.2.3. $\mathcal{H}_{s}=\mathcal{M}_{s}(\Omega, \mathcal{F}, \mu ; H) / \mathcal{J}_{s}(\Omega, \mathcal{F}, \mu ; H)$ will be called the Colombeau extension of $\mathbb{D}^{\infty}$, can also be called the space of Wiener-Colombeau distributions).
(Note. The subindex 's' is the reminiscent of the word 'simplified' . For $\mathbb{D}^{\infty}$ being replaced by $\mathcal{C}^{\infty}(\Omega) ; \Omega \subset \mathbb{R}^{n}$ and $\|\cdot\|_{k, p}$ norms by single index semi-norms in terms of supremums of partial derivatives on increasing compacts we retrieve the simplified Colombeau generalized functions, [16]).


Definition 2.2.5. Two members $\left(f_{\epsilon}\right)_{\epsilon},\left(g_{\epsilon}\right)_{\epsilon}$ of $\chi$ are said to be law equivalent if they have the same probability distribution.

## Conclusions.

1. $\mathcal{M}_{s}$ is an algebra: Let $\left(f_{\epsilon}\right)_{\epsilon},\left(g_{\epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$, define:
$\left(f_{\epsilon}\right)_{\epsilon} \cdot\left(g_{\epsilon}\right)_{\epsilon} \doteq\left(f_{\epsilon} \cdot g_{\epsilon}\right)_{\epsilon} \in \chi$ since $\mathbb{D}^{\infty} \equiv \mathbb{D}^{\infty}(\mathbb{R})$ is an algebra. Given $(k, p), k \in \mathbb{N}, 1<p<\infty, \exists N_{1}, N_{2} \in$ $\mathbb{N},\left\|f_{\epsilon}\right\|_{k, p}=O\left(\epsilon^{-N_{1}}\right)$ and $\left\|g_{\epsilon}\right\|_{k, p}=O\left(\epsilon^{-N_{2}}\right)$. Then $\left\|f_{\epsilon} . g_{\epsilon}\right\|_{[k, p}=O\left(\epsilon^{-N_{1}-N_{2}}\right)$, hence $\left(f_{\epsilon} . g_{\epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$.
2. $\mathcal{J}_{s}$ is an ideal in $\mathcal{M}_{s}:$ Let $\left(f_{\epsilon}\right)_{\epsilon} \in \mathcal{J}_{s}$ and $\left(g_{\epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$. Given $(k, p)$ as in $1 . \exists N_{1}, N_{2} \in \mathbb{N},\left\|f_{\epsilon}\right\|_{k, p}=O\left(\epsilon^{m-N_{1}}\right)$ $\forall m \geq N_{1},\left\|g_{\epsilon}\right\|_{k, p}=O\left(\epsilon^{-N_{2}}\right)$. Then $\left\|f_{\epsilon} . g_{\epsilon}\right\|_{k, p}=O\left(\epsilon^{m-N}\right) \forall m \geq N$, where $N=N_{1}+N_{2}$.
3. $\mathcal{H}_{s}$ is a factor algebra. If $F \in \mathcal{H}_{s}$, it is of the form $\left(f_{\epsilon}\right)_{\epsilon}+\mathcal{J}_{s},\left(f_{\epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$.
4. By the monotonicity of the norms $\|.\|_{k, p}$, if the conditions in Definitions 2.2.1 and 2.2.2 hold for a certain pair $(k, p)$, then they also hold for all pairs in the cone $\left\{\left(k^{\prime}, p^{\prime}\right): k^{\prime} \leq k, p^{\prime} \leq p\right\}$.

## Inclusions.

A) $\mathbb{D}^{\infty} \subset \mathcal{M}_{s}$ : If $F \in \mathbb{D}^{\infty}$, then take $\left(f_{\epsilon}\right)_{\epsilon}=F$ for all $\epsilon \in(0,1]$. Thus $\mathbb{D}^{\infty}$ is a faithful subalgebra of $\mathcal{H}_{s}$.
B) $\mathbb{D}^{-\infty} \subset \mathcal{H}_{s} \quad$ (inclusion of Meyer-Watanabe distributions)

Firstly consider the canonical $\mathbb{R}^{d}$-valued functional $W_{1}(\omega) \equiv \omega(1),(\omega \in \Omega$.) Then

$$
\begin{array}{r}
D W_{1, i}(t)=t e_{i} \in H, \quad D_{h} W_{1, i}(\omega)=h_{i}(1) \text { and thus } W_{1}(t) \in \mathbb{D}^{\infty}\left(\mathbb{R}^{d}\right) ; h \in H, t \in[0,1] \\
\left(e_{1}, \cdots, e_{d}\right) \quad \text { is an orthonormal basis in } \mathbb{R}^{d} \text { and }\left(D W_{1, i}, h\right)_{H}=h_{i}(1)=D_{h} W_{1, i} \tag{2.2.2}
\end{array}
$$

so that (2.1.2) is satisfied.

Given $T \in \mathbb{D}^{-\infty}$ we define the following functional in the convolution sense

$$
\begin{equation*}
f_{T, \epsilon}^{\phi}(\omega)=\left\langle T(\tilde{\omega}), \frac{1}{\epsilon^{d}} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\right\rangle, \epsilon \in(0,1], \omega, \tilde{\omega} \in \Omega ; \phi \in \mathcal{D}\left(\mathbb{R}^{d}\right) \tag{2.2.3}
\end{equation*}
$$

$\left(<.,>\right.$ denotes the bilinear form on $\left.\mathbb{D}^{\infty} \times \mathbb{D}^{-\infty}\right)$ and for fixed $\epsilon, \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right) \in \mathbb{D}^{\infty}$.

We show that, $f_{T, \epsilon}^{\phi} \in \chi$ and also it is moderate:

$$
\begin{gathered}
D_{h} f_{T, \epsilon}^{\phi}(\omega)=\lim _{\lambda \rightarrow 0} \frac{1}{\epsilon^{d}}\left\langle T\left(\tilde{\omega}, \phi\left(\frac{\omega(1)+\lambda h(1)-\tilde{\omega}(1)}{\epsilon}\right)-\phi\left(\frac{\omega(1)-\tilde{\omega}}{\epsilon}\right)\right\rangle\right. \\
\stackrel{(2.1 .2)}{=} \frac{1}{\epsilon^{d}}\left\langle T(\tilde{\omega}),\left(D \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right), h\right)_{H}\right\rangle \stackrel{(2.1 .4),(2.2 .1)}{=} \frac{1}{\epsilon^{d+1}}\left\langle T(\tilde{\omega}),\left(\sum_{i=1}^{d} \partial_{i} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right) t e_{i}, h\right)_{H}\right\rangle \\
=\frac{1}{\epsilon^{d+1}}\left\langle T(\tilde{\omega}), \sum_{i=1}^{d} \partial_{i} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right) \int_{0}^{1} \dot{h}_{i}(s) d s\right\rangle=\frac{1}{\epsilon^{d+1}} \int_{0}^{1} \sum_{i=1}^{d}\left\langle T\left(\tilde{\omega}, \partial_{i} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\right\rangle \dot{h}_{i}(s) d s\right. \\
=\frac{1}{\epsilon^{d+1}}\left(\sum_{i=1}^{d}<T(\tilde{\omega}), \partial_{i} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)>t e_{i}, h\right)_{H} .
\end{gathered}
$$

Hence by (2.1.2)

$$
\begin{equation*}
D f_{T, \epsilon}^{\phi}(\omega)=\frac{1}{\epsilon^{d+1}} \sum_{i=1}^{d}\left\langle T(\tilde{\omega}), \partial_{i} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\right\rangle t e_{i} \in L^{p}(H), t \in(0,1] \tag{2.2.4}
\end{equation*}
$$

(Note: (2.2.4) can be obtained informally from (2.2.3) by inserting the gradient in the right side of the bilinear form $<T(\tilde{\omega}), .>$, however this operation is not in general legitimate since the gradient of the right side is $H$-valued, thus does not belong to the domain of $T \in \mathbb{D}^{-\infty}$ ).

For the second gradient we use (2.1.3) with $E \longleftrightarrow H$ and for any $\tilde{h} \in H, \tau \in(0,1]$

$$
\begin{gather*}
\left(D_{h} D f_{T, \epsilon}^{\phi}, \tilde{h}\right)_{H}=\frac{1}{\epsilon^{d+2}}\left(\sum_{i=1}^{d}\left\langle T(\tilde{\omega}),\left(\sum_{j=1}^{d} \partial_{i, j}^{2} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right), \tau e_{j}, h\right)_{H}\right\rangle t e_{i}, \tilde{h}\right) \\
=\frac{1}{\epsilon^{d+2}}\left(\sum_{i=1}^{d}\left\langle T(\tilde{\omega}), \sum_{j=1}^{d} \partial_{i, j}^{2} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\left(\tau e_{j}, h\right)_{H}\right\rangle t e_{i}, \tilde{h}\right)_{H} \\
=\frac{1}{\epsilon^{d+2}} \sum_{i=1}^{d}\left\langle T(\tilde{\omega}) \sum_{j=1}^{d} \partial_{i, j}^{2} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\left(\tau e_{j}, h\right)_{H}\right\rangle\left(t e_{i}, \tilde{h}\right)_{H} \\
=\frac{1}{\epsilon^{d+2}} \sum_{i, j=1}^{d}\left\langle T(\tilde{\omega}), \partial_{i, j}^{2} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right\rangle\left(\tau e_{j}, h\right)_{H}\left(t e_{i}, \tilde{h}\right)_{H}\right. \\
=\frac{1}{\epsilon^{d+2}} \sum_{i, j=1}^{d}\left\langle T(\tilde{\omega}), \partial_{i, j}^{2} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\right\rangle\left(\tau e_{j} \otimes t e_{i}, h \otimes \tilde{h}\right)_{H \otimes H}=\left(D^{2} f_{T, \epsilon}^{\phi}, h \otimes \tilde{h}\right), \text { thus } \\
D^{2} f_{T, \epsilon}^{\phi}(\omega)=\frac{1}{\epsilon^{d+2}} \sum_{i, j=1}^{d}\left\langle T(\tilde{\omega}), \partial_{i, j}^{2} \phi\left(\frac{\omega(1)-\tilde{\omega}(2)}{\epsilon}\right)\right\rangle \tau e_{j} \otimes t e_{i} \in L^{p}(H \otimes H), \tau \in(0,1] \tag{2.2.5}
\end{gather*}
$$

By induction

$$
\begin{equation*}
D^{k} f_{T, \epsilon}^{\phi}=\frac{1}{d+k} \sum_{i_{1}, i_{2}, \cdots, i_{k}=1}^{k}\left\langle T(\tilde{\omega}), \partial_{i_{1}, i_{2}, \cdots, i_{k}} \phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\right\rangle \otimes_{r=1}^{k} t_{r} e_{i_{r}} ; \quad t_{r} \in(0,1],(r=1, \cdots, k) \tag{2.2.6}
\end{equation*}
$$

To verify the inductive hypothesis we should show by (2.1.3)

$$
\begin{equation*}
\left(D^{k+1} f_{T, \epsilon}^{\phi}, \otimes_{j=1}^{k+1} h_{j}\right)_{H^{\otimes k+1}}=\left(D_{h_{k+1}} f_{T, \epsilon}^{\phi}, \otimes_{j=1}^{k} h_{j}\right)_{H^{\otimes k}} ; h_{j} \in H,(j=1, \cdots, k+1) \tag{2.2.7}
\end{equation*}
$$

(The r,h.s. of (2.2.7) with the abbreviated symbol $\phi(\quad)$ )

$$
\begin{gathered}
\frac{1}{\epsilon^{d+k+1}}\left(\sum_{i_{1}, \cdots, i_{k}}^{d}<T(\tilde{\omega}), \sum_{i_{k+1}=1}^{d}\left(\partial_{i_{1}, \cdots, i_{k+1}}^{k+1} \phi(\quad) t_{k+1} e_{i_{k+1}}, h_{k+1}\right)_{H}>\otimes_{r=1}^{k} t_{r} e_{i_{r}}, \otimes_{j=1}^{k} h_{j}\right)_{H \otimes k} \\
=\frac{1}{\epsilon^{d+k+1}} \sum_{i_{1}, \cdots, i_{k+1}}^{d}<T(\tilde{\omega}), \partial_{i_{1}, \cdots, i_{k+1}}^{k+1} \phi(\quad)>\left(t_{k+1} e_{i_{k+1}}, h_{k+1}\right)_{H}\left(\otimes_{r=1}^{k} t_{r} e_{i_{r}}, \otimes_{j=1}^{k} h_{j}\right)_{H \otimes k} \\
\quad=\frac{1}{\epsilon^{d+k+1}}\left(\sum_{i_{1}, \cdots, i_{k+1}}^{d}<T(\tilde{\epsilon}), \partial_{i_{1}, \cdots, i_{k+1}} \phi(\quad) \otimes_{r=1}^{k+1} t_{r} e_{i_{r}}, \otimes_{j=1}^{k+1} h_{j}\right)_{H \otimes k+1}=\text { I.h.s. }
\end{gathered}
$$

With an easy estimation we find that $\left|D^{k} f_{T, \epsilon}^{\phi}\right|_{H \otimes k}^{p}<\infty$ so that $D^{k} f_{T, \epsilon}^{\phi} \in L^{p}\left(H^{\otimes k}\right), k \in \mathbb{N} /\{0\}$ Then referring to (2.1.5) we have $\left\|f_{T, \epsilon}^{\phi}\right\|_{k, p}<\infty$ and as ( $k, p$ ) are arbitrary, by (2.1.7) $\left(f_{T, \epsilon}^{\phi}\right)_{\epsilon} \in \chi$, furthermore $\left\|f_{T, \epsilon}^{\phi}\right\|_{k, p}=O\left(\epsilon^{-k-d}\right)$ hence $\left(f_{T, \epsilon}^{\phi}\right)_{\epsilon} \in \mathcal{M}_{s}$. Its class in $\mathcal{H}_{s}$ is a representative of $T \in \mathbb{D}^{-\infty}$.

If $\phi^{*} \in \mathcal{D}\left(\mathbb{R}^{d}\right)$ is another Schwartz test function, $f_{T, \epsilon}^{\phi}$ and $f_{T, \epsilon}^{\phi^{*}}$ are $\epsilon$-equivalent:
Recall that in our set-up the Donsker's delta function is given canonically by $\left.\delta_{x}(\omega)(1)\right)$ [19] where $\delta_{x}$ is the Dirac delta function at $x \in \mathbb{R}^{d}$. Then $\lim _{\epsilon \rightarrow 0}\left(f_{T, \epsilon}^{\phi}-f_{T, \epsilon}^{\phi^{*}}\right)=<T(\tilde{\omega}), \lim _{\epsilon \rightarrow 0}\left(\phi\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)-\phi^{*}\left(\frac{\omega(1)-\tilde{\omega}(1)}{\epsilon}\right)\right)>=$ $<T(\tilde{\omega}), \delta_{\tilde{\omega}(1)}-\delta_{\tilde{\omega}(1)}>=0$.

### 2.3 Sharp(uniform) Topology in $\mathcal{H}_{s}$

$\mathbb{D}^{\infty}$ is a countably normed space; for a strictly increasing set $\left\{p_{m}\right\},\left(1<p_{m}<\infty, m=1,2, \cdots\right)$ of numbers $\mathbb{D}^{\infty}=$ $\bigcap_{k \in \mathbb{N}} \bigcap \mathbb{D}_{k}^{p_{m}}$ (due to the monotonicity of $\|\cdot\|_{k, p}$ norms). Arrange the set of countable norms as

$$
\|\cdot\|_{1, p_{1}} ;\|\cdot\|_{2, p_{1}},\|\cdot\|_{1, p_{2}} ;\|\cdot\|_{3, p_{1}},\|\cdot\|_{2, p_{2}},\|\cdot\|_{1, p_{3}} ; \cdots \cdots ;\|\cdot\|_{n, p_{1}},\|\cdot\|_{n-1, p_{2}}, \cdots \cdots,\|\cdot,\|_{1, p_{n}} ; \cdots
$$

Then we define a new set of countable norms $\mu_{n},(n=1,2, \cdots$ by

$$
\begin{equation*}
\mu_{n}=\left(\sum_{j=1}^{n}\|\cdot\|_{n+1-j, p_{j}}^{2}\right)^{1 / 2} \tag{2.3.1}
\end{equation*}
$$

which are increasing; also $\left\{\mu_{n}\right\}_{n \in \mathbb{N}}$ and $\left\{\|\cdot\|_{k, p}\right\}$ norm systems are equivalent and create the same topology.
Furthermore $\left(f_{\epsilon}\right)_{\epsilon} \in \mathcal{H}_{s}$ is moderate in $\left\{\|\cdot\|_{k, p}\right\}$ norms $\Leftrightarrow$ it is moderate in $\left\{\mu_{n}\right\}$ norms.
For $\left(f_{\epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$ the $n$-valuation of $f$, denoted by $v_{n}(f)$ is $\sup _{b \in \mathbb{Z}}\left\{\mu_{n}\left(f_{\epsilon}\right)=O\left(\epsilon^{b}\right)\right\}$. A family $\delta_{n}$ of ultrametric pseudistances on $\mathcal{M}_{s}$ is defined by $\delta_{n}(f, g)=\exp \left(-v_{n}(f-g)\right), \forall f, g \in \mathcal{M}_{s}$.
$\left(f_{\epsilon}\right)_{\epsilon}$ lies in $\mathcal{J}_{s} \Longleftrightarrow \forall n \in \mathbb{N}, v_{n}(f)=+\infty$ (or equivalently $\forall n \in \mathbb{N}, \delta_{n}(f, 0)=0$ ).
These definitions transfer naturally to the quotient space $\mathcal{H}_{s}$. Valuations and pseudo-distances have properties well-known in analysis. The ultrametric uniform structure and the topology constructed are called sharp uniform structure and sharp topology respectively.
Theorem. The space $\mathcal{H}_{s}$ is complete for the sharp uniform structure.
Proof. (Based on [12], prop. 1.31 and [15], Thm 3.1). Let $\left(F_{n}\right)_{n \in \mathbb{N}}$ be a Cauchy sequence in $\mathcal{H}_{s}$ and let $\left(f_{\epsilon}\right)_{\epsilon}(n \in \mathbb{N})$ be a sequence of representatives in $\mathcal{M}_{s}$. That means $\forall n, \delta_{n}\left(f_{r}, f_{l}\right) \longrightarrow 0$ as $r \rightarrow \infty, l \rightarrow 0$. Considering the definition of the valuations we can extract a subsequence $\left(f_{q_{n}, \epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$ with strictly increasing sequence $q_{n}, n \in \mathbb{N}$ of integers having the following property:
$\exists$ a decreasing sequence $\epsilon_{k} \downarrow 0\left(\epsilon_{k} \leq \frac{1}{2^{k}}\right)$ such that

$$
\begin{equation*}
\forall \epsilon \in\left(0, \epsilon_{k}\right] \quad \mu_{k}\left(f_{q_{k}, \epsilon}-f_{q_{k-1}, \epsilon}\right)<\epsilon^{k}, \quad(k \geq 1) \tag{2.3.2}
\end{equation*}
$$

Define for $k \in \mathbb{N} /\{0\}$

$$
g_{k, \epsilon}= \begin{cases}f_{q_{k}, \epsilon}-f_{q_{k-1}, \epsilon} & \text { if } \epsilon \in\left(0, \epsilon_{k}\right)  \tag{2.3.3}\\ 0 & \text { if } \epsilon \in\left[\epsilon_{k}, 1\right]\end{cases}
$$

and finally define

$$
\begin{equation*}
f_{\epsilon} \doteq f_{q_{0}, \epsilon}+\sum_{k=1}^{\infty} g_{k, \epsilon} \tag{2.3.4}
\end{equation*}
$$

$f_{\epsilon}$ is locally finite. For any $n \in \mathbb{N} /\{0\}$

$$
\begin{gathered}
\mu_{n}\left(f_{\epsilon}\right) \leq \mu_{n}\left(f_{q_{0}, \epsilon}\right)+\sum_{k=1}^{\infty} \mu_{n}\left(g_{k, \epsilon}\right) \leq \mu_{n}\left(f_{q_{0}, \epsilon}\right)+\sum_{k=1}^{n-1} \mu_{n}\left(g_{k, \epsilon}\right)+\sum_{k=n}^{\infty} \mu_{k}\left(g_{k, \epsilon}\right) \\
\leq \mu_{n}\left(f_{q_{0}, \epsilon}\right)+\sum_{k=1}^{n-1} \mu_{n}\left(g_{k, \epsilon}\right)+\frac{\epsilon^{n}}{1-\epsilon}
\end{gathered}
$$

(since $\mu_{n}$ is increasing and by (2.3.2)). thus $\left(f_{\epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$
As $f_{\epsilon}=f_{q_{0}, \epsilon}+f_{q_{1}, \epsilon}-f_{q_{0}, \epsilon}+\cdots \cdots+f_{q_{m-1}, \epsilon}-f_{q_{m-2}, \epsilon}+f_{q_{m}, \epsilon}-f_{q_{m-1}, \epsilon}+\cdots$
We have by cancellations and for any $m \geq 1$

$$
f_{\epsilon}-f_{q_{m}, \epsilon}=\sum_{k=m}^{\infty} g_{k+1, \epsilon} ; \quad \mu_{m}\left(f_{\epsilon}-f_{q_{m}, \epsilon}\right)=\mu_{m}\left(\sum_{k=m}^{\infty} g_{k+1, \epsilon}\right) \leq \sum_{k=m}^{\infty} \mu_{k+1}\left(g_{k+1, \epsilon}\right) \leq \frac{\epsilon^{m+1}}{1-\epsilon}, \quad\left(\epsilon \in\left(0, \epsilon_{m}\right)\right)
$$

This shows that the subsequence $\left(f_{q_{n}}\right)$ converges to $\left(f_{\epsilon}\right)_{\epsilon} \in \mathcal{M}_{s}$ in sharp topology.
Thus the subsequence $\left(F_{q_{n}}\right)$ converges to $F \in \mathcal{H}_{s},\left(F_{q_{n}}\right.$ and $F$ are the classes of $\left(f_{q_{n}, \epsilon}\right)_{\epsilon}$ and $\left(f_{\epsilon}\right)_{\epsilon}$ in $\mathcal{H}_{s}$ respectively). As $\left(F_{n}\right)_{n \in \mathbb{N}}$ is Cauchy with a convergent subsequence it converges to $F$. $\square$
Note. $\mathcal{H}_{s}$ is also a $\overline{\mathbb{C}}$ topological module, $\overline{\mathbb{C}}$ being the ring of generalized complex numbers. But this result is not needed in the sequel.

## 3 Inclusion of Hida Distributions

### 3.1 Real Case

We consider the Gelfand triplet $E \subset H \subset E^{*} ; H$ is a real separable Hilbert space, $E$ is a countably Hilbert nuclear space which is densely and continuously imbedded in $H$ and $E^{*}$ is the dual of $E$.
By Minlos theorem there is a unique Gaussian measure $\mu$ on $\left(E^{*}, \mathcal{B}\left(E^{*}\right)\right.$ ). Thus ( $\left.E^{*}, \mathcal{B}\left(E^{*}\right), \mu ; H\right)$ forms a canonical Gaussian probability space. We denote the inner product and norm in $H\left(\right.$ or $\left(H^{\otimes n}\right)$ by (.,.) and |.| respectively. $L^{2}\left(E^{*}, \mu\right)$ is abbreviated by ( $L^{2}$ ).
Using the second quantization operator we construct the Hida testing functional space $(E)$ and its dual $(E)^{*}$, the Hida distributions.
On the other hand the Meyer-Watanabe functional space $\mathbb{D}^{\infty}$ and the distribution space $\mathbb{D}^{-\infty}$ can also be constructed on $\left(E^{*}, \mathcal{B}\left(E^{*}\right), \mu\right)$ as well, the C-M space being replaced by $H$.
The starting point of the weak differential calculus will be then

$$
\begin{equation*}
(D F, h)_{H}=\lim _{\lambda \rightarrow 0} \frac{F(x+\lambda h)-F(x)}{\lambda}, x \in E^{*}, h \in H \tag{3.1.1}
\end{equation*}
$$

For $f \in E$, consider the functional $W_{f}$, defined by

$$
\begin{equation*}
W_{f}(x) \doteq<x, f>, \quad\left(<., .>\text { denotes the bilinear form on } E^{*} \times E \text { or on } E \times E^{*}\right) \tag{3.1.2}
\end{equation*}
$$

By Minlos' theorem $\int_{E^{*}} e^{i<x, f>} d \mu(x)=e^{-\frac{1}{2}|f|^{2}}$, hence $W_{f}$ is a Gaussian random variable on $E^{*}$.
The linear map $f \longrightarrow W_{f}: E \longrightarrow L^{2}\left(E^{*}, \mu\right)$ can be extended to a linear isometry from $H$ into $L^{2}\left(E^{*}, \mu\right)$. Now

$$
\begin{equation*}
\left(D W_{f}, h\right)_{H}=\lim _{\lambda \rightarrow 0} \frac{<f, x+\lambda h>-<f, x>}{\lambda}=<f, h>=(f, h)_{H} \quad \text { yielding } \quad D W_{f}=f \in E \subset H \tag{3.1.3}
\end{equation*}
$$

It is known that for $\left(\mathbb{D}^{\infty}, \mathbb{D}^{-\infty}\right)$ constructed on $\left(E^{*}, \mathcal{B}\left(E^{*}\right), \mu\right)$ we have $(E) \subset \mathbb{D}^{\infty}$ and $\mathbb{D}^{-\infty} \subset(E)^{*}$ both inclusions being dense. For $T \in(E)^{*}$ the counterpart of (2.2.3) will be $(d=1)$

$$
\begin{equation*}
g_{T, \epsilon}^{\phi, f}(x)=\frac{1}{\epsilon} \ll T(\tilde{x}), \phi\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right) \gg ; x, \tilde{x} \in E^{*}, \phi \in \mathcal{D}(\mathbb{R}),|f|_{H}=1 \tag{3.1.4}
\end{equation*}
$$

where $\ll ., . \gg$ is the canonical bilinear form on $\left((E)^{*},(E)\right)$.

Polynomial functionals in the variables $\left\{W_{f_{j}}\right\}(j=1,2, \cdots n)$ belong to $\mathbb{D}^{\infty} \cap(E)$ and are dense in both of them. The right member of $\ll ., . \gg$ in (3.1.4) can be regarded as the limit of Taylor polynomials in $W_{f}$ and is in ( $E$ ) due to its completeness as a countably Hilbertian nuclear space (in fact a smooth functional).
We show that $\left(g_{T, \epsilon}^{\phi, f}\right)_{\epsilon} \in \chi$ and that it is also moderate. We have:

$$
\begin{gathered}
\left(D g_{\epsilon}, h\right)_{H}=\frac{1}{\epsilon} \ll T(\tilde{x}),\left(D \phi\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right), h\right)_{H} \gg=\frac{1}{\epsilon^{2}} \ll T(\tilde{x}),\left(\phi^{\prime}\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right) f, h\right)_{H} \gg \\
=\frac{1}{\epsilon^{2}} \ll T(\tilde{x}), \phi^{\prime}\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right) \gg(f, h)_{H}=\frac{1}{\epsilon^{2}}\left(\ll T(\tilde{x}), \phi^{\prime}\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right) \gg f, h\right)_{H}
\end{gathered}
$$

which yields $D g_{\epsilon}=\frac{1}{\epsilon^{2}} \ll T(\tilde{x}),\left(\phi^{\prime}\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right) f \gg\right.$.
For higher order gradients we again use (2.1.3), in some abbreviated notation :

$$
\begin{gathered}
\left(D_{h} D g_{\epsilon}, \tilde{h}\right)_{H}=\frac{1}{\epsilon^{3}}\left(\ll T(\tilde{x}),\left(\phi^{\prime \prime}() f, h\right)_{H} \gg f, \tilde{h}\right)_{H}=\frac{1}{\epsilon^{3}}\left(\ll T(\tilde{x}), \phi^{\prime \prime}()(f, h)_{H} \gg f, \tilde{h}\right)_{H} \\
=\frac{1}{\epsilon^{3}} \ll T(\tilde{x}), \phi^{\prime \prime}(\quad) \gg(f, h)_{H}(f, \tilde{h})_{H}=\frac{1}{\epsilon^{3}} \ll T\left(\tilde{x}, \phi^{\prime \prime}() \gg(f \otimes f), h \otimes \tilde{h}\right)_{H \otimes H} \\
=\frac{1}{\epsilon^{3}}\left(\ll T(\tilde{x}), \phi^{\prime \prime}() \gg(f \otimes f, h \otimes \tilde{h})_{H \otimes H}=\left(D^{2} g_{\epsilon}, h \otimes \tilde{h}\right)_{H \otimes H}\right.
\end{gathered}
$$

giving $D^{2} g_{T, \epsilon}^{\phi, f}(x)=\frac{1}{\epsilon^{3}} \ll T(\tilde{x}),, \phi^{\prime \prime}\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right) \gg(f \otimes f)$.
By an induction similar to the one in Section 2.2:

$$
\begin{equation*}
D^{k} g_{T, \epsilon}^{\phi, f}(x)=\frac{1}{\epsilon^{k+1}} \ll T(\tilde{x}), \phi^{(k)}\left(\frac{W_{f}(x)-W_{f}(\tilde{x})}{\epsilon}\right) \gg f^{\otimes k}=\frac{1}{\epsilon^{k+1}} \ll T(\tilde{x}), \phi^{(k)}\left(\frac{<x, f>-<\tilde{x}, f>}{\epsilon}\right) \gg f^{\otimes k} \tag{3.1.5}
\end{equation*}
$$

As clearly $\left|D^{k} g_{T, \epsilon}^{\phi, f}\right|_{H \otimes k}^{p}<\infty$, we have $D^{k} g_{T, \epsilon}^{\phi, f} \in L^{p}\left(H^{\otimes k}\right)$. so that by (2.1.5) $\left(g_{T, \epsilon}^{\phi, f}\right)_{\epsilon} \in \chi$ and furthermore $\left(g_{T, \epsilon}^{\phi, f}\right)_{\epsilon} \in$ $\mathcal{M}_{s}$ as $\left\|g_{T, \epsilon}^{\phi, f}\right\|_{k, p}=O\left(\epsilon^{-k-1}\right)$, hence $(E)^{*} \in \mathcal{H}_{s}$.
Different representatives obtained by $\tilde{\phi} \in \mathcal{D}(\mathbb{R})$ and/or by $\tilde{f} \in E,|\tilde{f}|_{H}=1$, are either $\epsilon$-equivalent or law equivalent or a combination of them according to Definitions 2.2.4, 2.2.5.

### 3.2 Complex case

For any topological vector space $K$ on $\mathbb{R}$, denote by $V_{\mathbb{C}}$ its complexification, i.e. $V_{\mathbb{C}}=V+i V$.
If $V$ is a Hilbert space and $u_{1}+i v_{1}$ and $u_{2}+i v_{2}$, are in $V_{\mathbb{C}}$, then their inner product is, $\left(u_{1}+i v_{1}, u_{2}+i v_{2}\right)_{V_{\mathbb{C}}}=$ $\left(u_{1}, u_{2}\right)+\left(v_{1}, v_{2}\right)+i\left[\left(u_{2}, v_{1}\right)-\left(u_{1}, v_{2}\right)\right]$ thus $\left|u_{1}+i v_{1}\right|_{V_{\mathrm{C}}}^{2}=\left|u_{1}\right|_{V}^{2}+\left|v_{1}\right|_{V}^{2}$.
Let $\chi_{\mathbb{C}}$ be the set of all functions $(0,1] \longrightarrow \mathbb{D}_{\mathbb{C}}^{\infty}=\mathbb{D}^{\infty}+i \mathbb{D}^{\infty}$. Definitions 2.2.1 to 2.2.5 are modified accordingly to define $\mathcal{H}_{\mathbb{C}, s}, \epsilon-$ and law equivalence.
$\mathcal{H}_{\mathbb{C}, s}$ is also an algebra. We consider complex-valued functionals on the measure space $\left(E^{*}, \mathcal{B}\left(E^{*}\right), \mu\right)$. Letting $E_{\mathbb{C}} \ni f, f=$ $f_{1}+i f_{2},\left(f_{j} \in E, j=1.2\right)$
For $x \in E^{*}, W_{f}=<x, f_{1}+i f_{2}>=<x, f_{1}>+i<x, f_{2}>$.
The real and imaginary parts of $W_{f}$ are independent Gaussian random variables with characteristic functions $e^{-\frac{1}{2}\left|f_{j}\right|^{2}},(j=$ $1,2)$. If we take $\left|f_{1}\right|^{2}=\left|f_{2}\right|^{2}=1 / 2$, then $W_{f} \in \mathbb{C} N(0,1)$, i.e. a standard complex-valued Gaussian random variable.
The linear map $f \rightarrow W_{f}$ from $E_{\mathbb{C}}$ to $L_{\mathbb{C}}^{2}\left(E^{*}, \mu\right)=\left(L^{2}\right)_{\mathbb{C}}$ can be extended to a linear isometry $H_{\mathbb{C}} \rightarrow\left(L^{2}\right)_{\mathbb{C}}$.
Consider the formal composition $\delta\left(W_{f}(x)-W_{f}(\tilde{x})\right), x, \tilde{x} \in E^{*}$, where $\delta$ is the Dirac delta function. Using the linearity of Dirac delta it takes the form $\delta\left(<x, f_{1}>-<\tilde{x}, f_{1}>\right)+i \delta\left(<x, f_{2}>-<\tilde{x}, f_{2}>\right)$. This can be approximateed by the delta nets as
$\frac{1}{\epsilon} \phi\left(\frac{\leq x, f_{1}>-<\tilde{x}, f_{1}>}{\epsilon}\right)+i \frac{1}{\epsilon} \phi\left(\frac{\left\langle x, f_{2}>-<\tilde{x}, f_{2}>\right.}{\epsilon}\right), \phi \in \mathcal{D}(\mathbb{R}) \quad(\dagger)$
Its limit can be regarded as a complex linear combination of two functions of so-called Donsker's type, (see also the next section ). Let $T \in(E)_{\mathbb{C}}^{*}$, based on ( $\dagger$ ) define

$$
\begin{equation*}
\Theta_{T, \epsilon}^{\phi, f}(x)=\frac{1}{\epsilon} \ll T(\tilde{x}), \phi\left(\frac{<x, f_{1}>-<\tilde{x}, f_{1}>}{\epsilon}\right)+i \phi\left(\frac{<x, f_{2}>-<\tilde{x} f_{2}>}{\epsilon}\right) \gg, f=f_{1}+i f_{2}, x, \tilde{x} \in E^{*} \tag{3.2.1}
\end{equation*}
$$

where $\ll ., . \gg$ is the canonical conjugate bilinear form on $(E)_{\mathbb{C}}^{*} \times\left(E_{\mathbb{C}}\right)$. The use of two tests functions $\phi$ will not bring in an essential difference.
The real and imaginary parts of the right side of $\ll ., . \gg$ are both in $(E)$ according to Section 3.1. Therefore the right side belongs to $(E)_{\mathbb{C}}$. (For a complete characterization of $(E)_{\mathbb{C}}$ and $(E)_{\mathbb{C}}^{*}$ in terms of S-transforms and $U_{\beta}$ functionals, c.f [19] and [21]).

The calculation of the gradients of the functional (3.2.1) is basically the same as in the previous section except obvious modifications for the complex case. Thus

$$
\begin{equation*}
D \Theta_{T, \epsilon}^{\phi, f}(x)=\frac{1}{\epsilon^{2}} \ll T(\tilde{x}), \phi^{\prime}\left(\frac{<x, f_{1}>-<\tilde{x}, f_{1}>}{\epsilon}\right) \gg f_{1}+i \frac{1}{\epsilon^{2}} \ll T(\tilde{x}), \phi^{\prime}\left(\frac{<x, f_{2}>-<\tilde{x}, f_{2}>}{\epsilon}\right) \gg f_{2} \tag{3.2.2}
\end{equation*}
$$

and

$$
\begin{gather*}
D^{k} \Theta_{T, \epsilon}^{\phi, f}(x)=\frac{1}{\epsilon^{k+1}} \ll T(\tilde{x}), \phi^{(k)}\left(\frac{<x, f_{1}>-<\tilde{x}, f_{1}>}{\epsilon}\right) \gg f_{1}^{\otimes k} \\
+i \frac{1}{\epsilon^{k+1}} \ll T(\tilde{x}), \phi^{(k)}\left(\frac{<x, f_{2}>-<\tilde{x}, f_{2}>}{\epsilon}\right) \gg f_{2}^{\otimes k} \tag{3.2.3}
\end{gather*}
$$

In parallel to the evaluations of Section 3.1, the $\|\cdot\|_{k, p}$ norms of the real and imaginary parts of the right-hand sides in this expression are finite. This then shows that for arbitrary $k \in \mathbb{N}$ and $1<p<\infty$ we have $\left\|\Theta_{T, \epsilon}^{\phi, f}\right\|_{k, p}<\infty$. Thus $\Theta_{T, \epsilon}^{\phi, f} \in \chi_{\mathbb{C}}$ and it is also moderate and represents $T \in(E)_{\mathbb{C}}^{*}$ in $\mathcal{H}_{\mathbb{C}, s}$. In this way $(E)_{\mathbb{C}}^{*} \subset \mathcal{H}_{\mathbb{C}, s}$.

### 3.3 Application to the Feynman Integrand.

Recall that Donsker's delta function in white noise theory set-up is given by $\delta_{a}\left(W_{f}\right) \equiv \delta\left(W_{f}-a\right)$
$; W_{f}(x)=<x, f>, x \in E^{*}, f \in E, a \in \mathbb{R}$.
It is shown to be a Hida distribution, therefore included in $\mathcal{H}_{s}$ with a representative $R(x)=\frac{1}{\epsilon} \phi\left(\frac{W_{f}(x)-a}{\epsilon}\right), \phi \in \mathcal{D}(\mathbb{R})$. As in section 3.1 we obtain

$$
\begin{equation*}
D R(x)=\frac{1}{\epsilon^{2}} \phi^{\prime}\left(\frac{<f, x>-a}{\epsilon}\right) f, \cdots, D^{k} R(x)=\frac{1}{\epsilon^{k+1}} \phi^{(k)}\left(\frac{<f, x>-a}{\epsilon}\right) f^{\otimes k} \tag{3.3.1}
\end{equation*}
$$

The other representatives are $\epsilon$-equivalent.
$\delta\left(B_{t}\right)$, where $B_{t}$ is the standard Brownian motion starting from 0 is also a Donsker's delta function. In the representative we may select $|f|=1, f \in E$ and use $W_{\sqrt{t} f}$.
$\mathcal{E}_{f}=\exp \left\{W_{f}-\frac{1}{2}<f, f>\right\}$ is the exponential functional and for $f \in E_{\mathbb{C}}$ it is known to be in $(E)_{\mathbb{C}}$.

For all $f \in E$, and $\lambda>0, G(f)=\int_{E^{*}} \mathcal{E}_{f} \mu^{(\lambda)}(d x)=\exp \left\{\frac{\lambda^{2}-1}{2}|f|^{2}\right\}$ is shown to be the $S$-transform of a Hida distribution $F(\lambda)=S^{-1} G$, [19].
For $\lambda \in \mathbb{C}, G(f)$ is still the $S$-transform of some element $F \in(E)_{\mathbb{C}}^{*}$, i.e. $F(\lambda)=S^{-1} G$.
The following short review is from [19] , [22]:
(In this section the initial Gelfand triplet of Section 3.1 can be taken as $S\left(\mathbb{R}^{d}\right) \subset L^{2}\left(\mathbb{R}^{d}\right) \subset S^{*}\left(\mathbb{R}^{d}\right)$ where $S^{*}\left(\mathbb{R}^{d}\right)$ denotes the set of tempered distributions ).
The Schrödinger equation

$$
\begin{equation*}
\frac{\partial \Psi}{\partial t}=i\left(\frac{\Delta}{2}-V\right) \Psi ; \quad \Psi(0, x)=f(x) \tag{3.3.2}
\end{equation*}
$$

where $\Delta$ is the Laplace operator in $\mathbb{R}^{d}, V$ is a real Borel function on $\mathbb{R}^{d}$, (the potential). From the point of white-noise analysis it is preferable to start with the heat equation

$$
\begin{equation*}
\frac{\partial u}{\partial t}=\left(\frac{\lambda}{2} \Delta-i V\right) u ; u(0, x)=f(x), \text { where } \lambda>0 \tag{3.3.3}
\end{equation*}
$$

The solution of (3.3.3) is given by the Feynman-Kac formula as

$$
\begin{equation*}
u(t, x)=\mathbb{E}\left[f\left(\sqrt{\lambda} B_{t}+x\right) \exp -i \int_{0}^{t} V(\sqrt{\lambda} s+x d s]\right. \tag{3.3.4}
\end{equation*}
$$

where $B_{t}$ is the standard Brownian motion starting from 0. In term of the Hida distribution $F(\sqrt{\lambda})$ this can be rewritten as

$$
\begin{equation*}
u(t, x)=\ll F(\sqrt{\lambda}) f\left(B_{t}+x\right) \exp -i \int_{0}^{t} V\left(B_{s}+x\right) d x, 1 \gg \tag{3.3.5}
\end{equation*}
$$

Let $f(x)=\delta_{y}(x)$. Then by (3.3.4) the fundamental solution of (3.3.3) is obtained as

$$
\begin{equation*}
u^{\lambda}(t, x, y)=\ll F(\sqrt{\lambda}) \delta\left(B_{t}-y+z\right) \exp -i \int_{0}^{t} V\left(B_{s}+x\right) d x, 1 \gg \tag{3.3.6}
\end{equation*}
$$

Suppose $u^{\lambda}$ has an analytic continuation in $\lambda$, (the conditions for validity of this analytic continuation are given at Corollary 4.4 [23] ), then the fundamental solution of (3.3.2) can be asserted to be

$$
\begin{equation*}
\Psi^{\lambda}(t, x, y)=\ll F(\sqrt{i}) \delta\left(B_{t}-y+x\right) e^{-i \int_{0}^{t} V\left(B_{s}+x\right) d s}, 1 \gg ; F(\sqrt{i}) \in(E)_{\mathbb{C}}^{*} \tag{3.3.7}
\end{equation*}
$$

The two factors $F(\sqrt{i})$ and $\delta\left(B_{t}-y+z\right)$ are in $\mathcal{H}_{s}$ by Section 3.2 and the paragraphs above. Therefore it remains to discuss the meaning of the exponential factor. $V\left(B_{s}+x\right)$ is expressed via a Bochner integral as

$$
\begin{equation*}
V\left(B_{s}+x\right)=\int_{\mathbb{R}^{d}} V(x) \delta\left(B_{s}+x-z\right) d z \tag{3.3.8}
\end{equation*}
$$

For $\Delta_{n}=\left\{\left(t_{1}, \cdots, t_{n}\right): 0<t_{1}<\cdots<t_{n}<t\right\}$

$$
\begin{align*}
& \exp \left\{-i \int_{0}^{t} V\left(B_{s}+x\right) d s\right\}=\sum_{n=0}^{\infty}(-i)^{n} \int_{\Delta_{n}} \prod_{j=1}^{n} V\left(B_{t_{j}}+x\right) d^{n} t \\
&=\sum_{n=0}^{\infty}(-i)^{n} \int_{\Delta_{n}} d t^{n} \prod_{j=1}^{n} \int_{\mathbb{R}^{d}} V\left(z_{j}\right) \delta\left(B_{t_{j}}+x-z_{j}\right) d z_{j} \tag{3.3.9}
\end{align*}
$$

so that the full Feynman integrand for the propagator becomes

$$
\begin{gather*}
\sum_{n=0}^{\infty}(-i)^{n} \int_{\Delta_{n}} d^{n} t \prod_{j=1}^{a} r e n \int_{\mathbb{R}^{d}} V\left(z_{j}\right) \delta\left(B_{t_{j}}+x-z_{j}\right) d z_{j} F(\sqrt{i}) \delta\left(B_{t}-y+x\right)=\sum_{n=0}^{\infty}(-i)^{n} \int_{\Delta_{n}} d^{n} t \prod_{j=1}^{n} \int_{\mathbb{R}^{d}} V\left(z_{j}\right) d z_{j} \\
\times F(\sqrt{i}) \delta\left(B_{t}-y+x\right) \prod_{j=1}^{n} \delta\left(B_{t_{j}}+x-z_{j}\right) \tag{3.3.10}
\end{gather*}
$$

In [22], [23], [24 ]considerable effort is spent to give a reasonable meaning to the factor $F(\sqrt{i}) \delta(\quad)$ as a Hida distribution. But obviously this factor is in algebra $\mathcal{H}_{s, \mathrm{C}}$. furthermore in [21] it is shown that under the assumptions that $d=1$, and $V(y) d y$ is a compactly supported signed measure, the product of the three terms in the right-hand side of (3.3.5) is a Hida distribution. To perform the $t_{i}$-integrations in (3.3.9) and the convergence of the series in the first factor of (3.3.9) the assumption $d=1$ is essential. For in the evaluation the integral $M_{n}=\int_{\Delta_{n}} d^{n} t \prod\left|t_{i}-t_{i-1}\right|^{-\frac{d}{2}}$ is needed and it exists only for $d=1$. In this case $M_{n}$ rapidly decreases and $\sum_{n=1}^{\infty} M_{n}<\infty$.
However in our case when $d=1$, without needing the introduction of $\Delta_{n}$ and the elaborate evaluations of (3.3.8) and (3.3.9), we can show the following:

Theorem. If $d=1$ and $V(z) d z$ is a bounded signed measure, then the product of three terms in (3.3.6) is in $\mathcal{H}_{s, \mathrm{C}}$.

Proof. Firstly let us show that $V\left(B_{s}+x\right)=\int_{\mathbb{R}} V(z) \delta\left(B_{s}+x-z\right) d z$ given by (3.3.7) is in $\mathcal{H}_{s}$. It will have a representative

$$
\begin{equation*}
\Gamma_{\epsilon}^{\phi, f}(\psi)=\frac{1}{\epsilon} \int_{\mathbb{R}} V(z) \phi\left(\frac{\left(W_{\sqrt{s} f}(\psi)-(z-x)\right.}{\epsilon}\right) d z, \quad f \in E,|f|=1, \phi \in \mathcal{D}(\mathbb{R}), \psi \in E^{*} \tag{3.3.11}
\end{equation*}
$$

By bounded convergence theorem the operators $D_{h}$ and $D$, can be inserted into the integral and an approach parallel to those in sections 2.2 and 3.1 yields

$$
D \Gamma_{\epsilon}^{\phi, f}(\psi)=\frac{1}{\epsilon^{2}}\left[\int_{\mathbb{R}} \phi^{\prime}\left(\frac{\sqrt{s}<\psi, x>-(z-x)}{\epsilon}\right) \sqrt{s} V(z) d z\right] f
$$

and

$$
\begin{gather*}
D^{k} \Gamma_{\epsilon}^{\phi, f}(\psi)=\frac{1}{\epsilon^{x+1}}\left[\int_{\mathbb{R}} \phi^{(k)}\left(\frac{\sqrt{s}<\psi, f>-(z-x)}{\epsilon}\right) s^{k / 2} V(z) d z\right] f^{\otimes k}  \tag{3.3.12}\\
\left|D^{k} \Gamma_{\epsilon}^{\phi, f}\right|_{H}^{p} \otimes k=\frac{1}{\epsilon^{k+1}}\left|\int_{\text {supp } \phi^{(k)}} \phi^{(k)}\left(\frac{\sqrt{s}<\psi, f>-(z-x)}{\epsilon}\right) s^{k / 2} V(z) d z\right|^{p}<\infty \tag{3.3.13}
\end{gather*}
$$

since $\left|f^{\otimes k}\right|_{H^{\otimes k}}=1$. As $V(z) d z=V(z)^{+} d z-V(z)^{-} d z$ is a bounded signed measure (3.3.12) shows that $\left\|D^{k} \Gamma_{\epsilon}^{\phi, f}\right\|_{L^{p}\left(\mu ; H^{\otimes k}\right)}<$ $\infty$ As $k \in \mathbb{N}$ and $1<p<\infty$ are arbitrary it follows from (2.1.5) that $\Gamma_{\epsilon}^{\phi, f} \in \mathbb{D}^{\infty}$, ( $0<\epsilon<1$ ), furthermore $\left\|\Gamma_{\epsilon}^{\phi, f}\right\|_{k, p}=O\left(\epsilon^{-k-1}\right)$ hence $\left(\Gamma_{\epsilon}^{\phi, f}\right)_{\epsilon} \in \mathcal{M}_{s}$, then its class in $\mathcal{H}_{s}$ represents $V\left(B_{s}+x\right)$.
$\int_{0}^{t} V\left(B_{s}+x\right) d s$ regarded as the limit of finite sums over the partitions of $[0,1]$ is in $\mathcal{H}_{s}$ due to its completion.
Now $e^{-i \int_{0}^{t} V\left(B_{s}+x\right) d s}=\cos \int_{0}^{t} V\left(B_{s}+x\right) d s-i \sin \int_{0}^{t} V\left(B_{s}+x\right)$, and as a complex linear combination of smooth, bounded functions of members of $\mathcal{M}_{s}$, (2.1.4) this expression is included in $\mathcal{M}_{s, \mathrm{C}}$. Its class in $\mathcal{H}_{s, \mathrm{C}}$ represents $e^{-i \int_{0}^{t} V\left(B_{s}+x\right) d s}$. As a consequence each factor of the product in (3.3.6) is in $\mathcal{H}_{s, \mathrm{C}}$, so is their product.

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## Concluding Remarks.

1) Present approach may have potentiality in studying the generalized solutions of non-linear stochastic partial differential equations as in [11].
2) In this paper non-linearities of polynomial order of $1 / \epsilon$ was considered. For non-linearities growing faster than polynomial order (e.g. of exponential order) an asympotic scale $a_{n}(\epsilon)$ accompanied by a set of postulates can be introduced, (c.f. [16]). In our case $a_{n}(\epsilon)=\epsilon^{n}$.

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